



***Assessing massive black hole
cosmic history through
gravitational waves***

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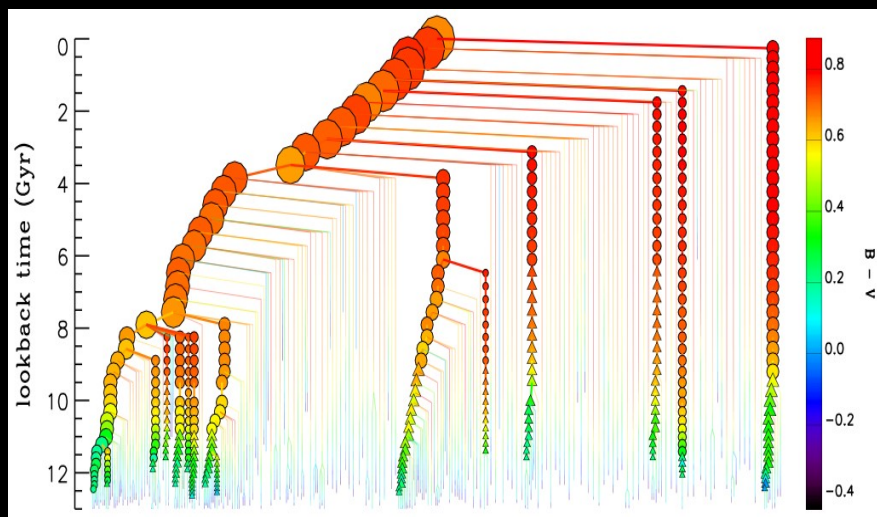
In collaboration with:

Jonathan Gair

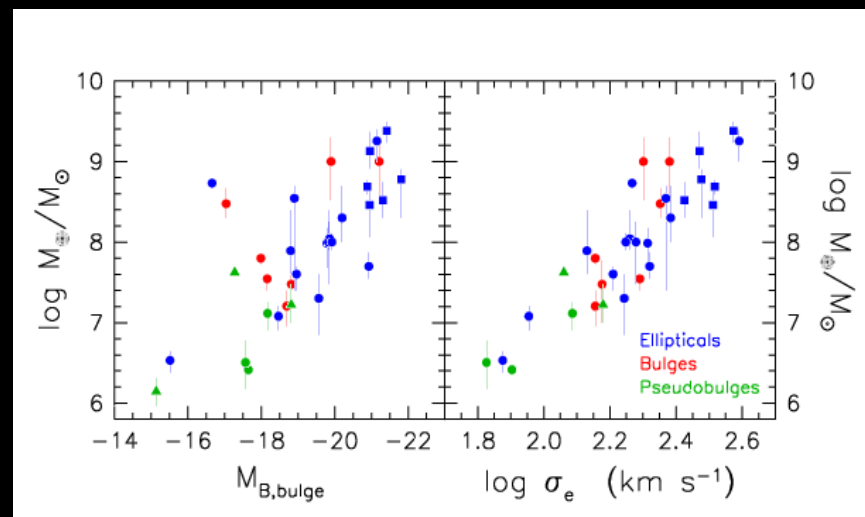
Emanuele Berti

Marta Volonteri

MBHBs are a natural consequence of the hierarchical formation of cosmic structures

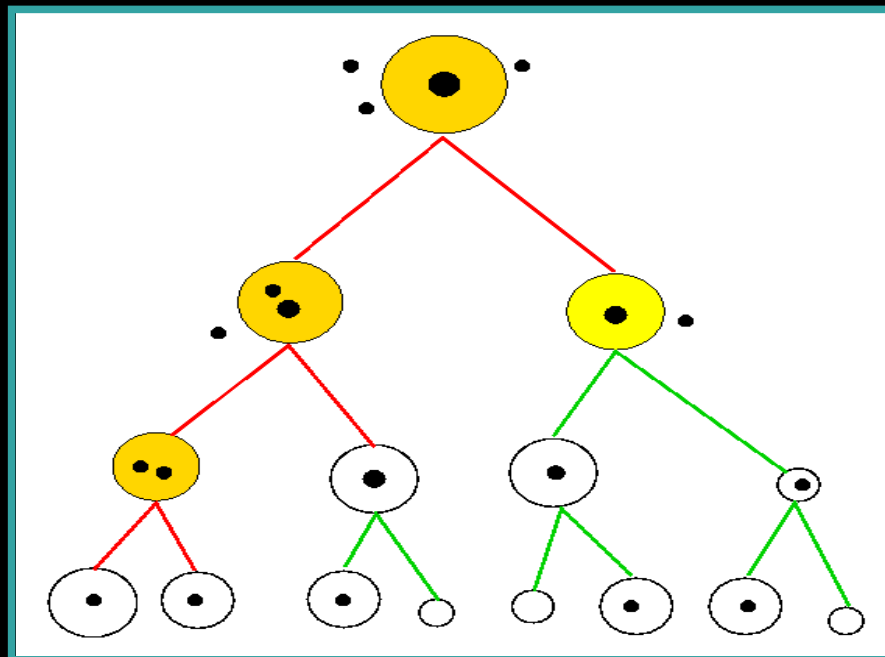


From De Lucia et al 2006



Ferrarese & Merritt 2000, Gebhardt et al. 2000

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Volonteri Haardt & Madau 2003

Question:

Given a set of observed massive black hole coalescences,
What information can be extracted about the underlying massive black
hole population model?

We consider 10 different formation models differing in:

- 1- MBH seeding mechanism (heavy vs light seeds)
- 2- Metallicity feedback (metal free vs all metallicities)
- 3- Accretion efficiency (Eddington vs. Merloni & Heinz)
- 4- Accretion geometry (coherent vs. chaotic)

Name	i	Seeding	Metallicity	Accretion model	Accretion geometry	\bar{N}_i [yr^{-1}]
<i>VHM-noZ-Edd-co</i>	1	POPIII	$Z = 0$	Eddington	coherent	86
<i>VHM-noZ-Edd-ch</i>	2	POPIII	$Z = 0$	Eddington	chaotic	81
<i>VHM-Z-Edd-co</i>	3	POPIII	all Z	Eddington	coherent	108
<i>VHM-Z-Edd-ch</i>	4	POPIII	all Z	Eddington	chaotic	113
<i>BVR-noZ-Edd-co</i>	5	Quasistar	$Z = 0$	Eddington	coherent	26
<i>BVR-noZ-Edd-ch</i>	6	Quasistar	$Z = 0$	Eddington	chaotic	24
<i>BVR-Z-Edd-co</i>	7	Quasistar	all Z	Eddington	coherent	22
<i>BVR-Z-Edd-ch</i>	8	Quasistar	all Z	Eddington	chaotic	29
<i>BVR-noZ-MH-co</i>	9	Quasistar	$Z = 0$	Merloni & Heinz	coherent	33
<i>BVR-noZ-MH-ch</i>	10	Quasistar	$Z = 0$	Merloni & Heinz	chaotic	33

PROCEDURE

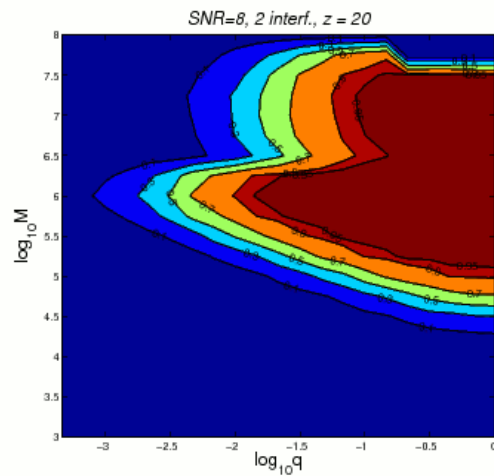
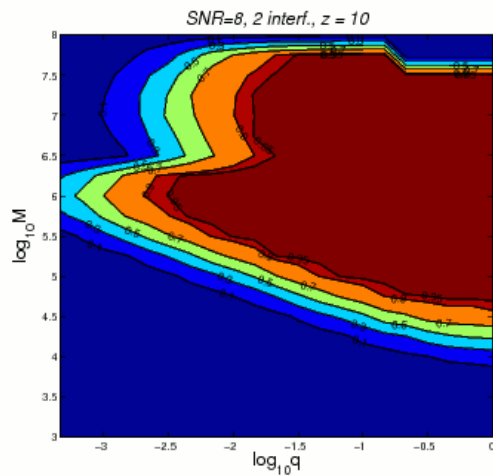
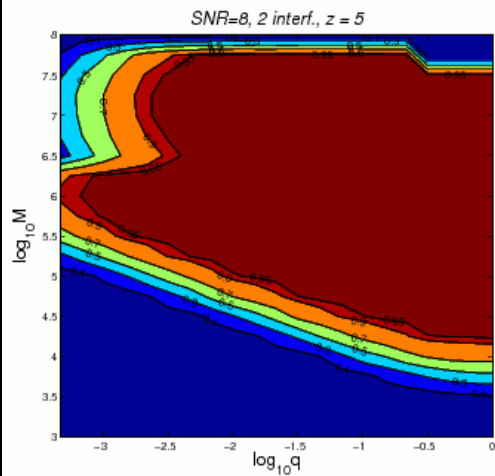
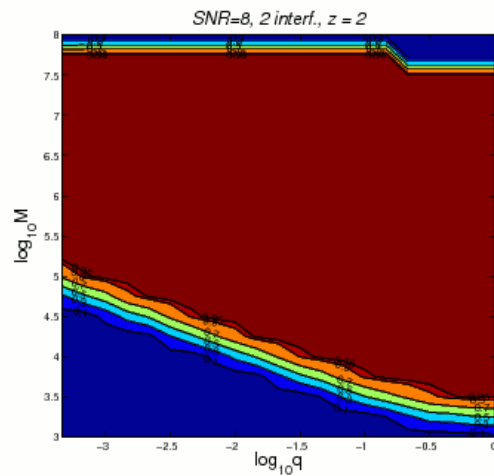
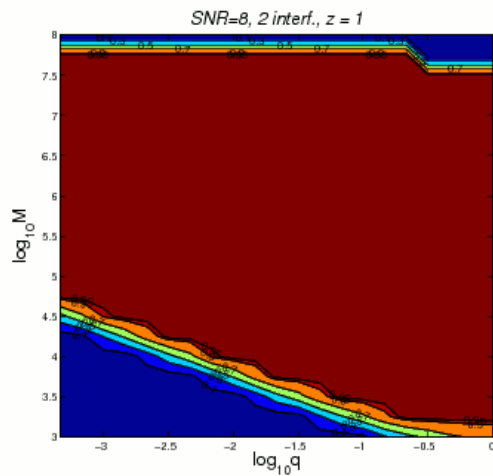
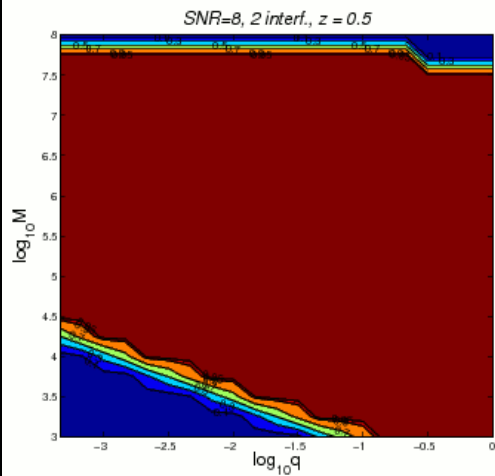
- a-Construct the detector transfer function (takes into account for the adopted waveform and for the detector performance)
- b-Filter the theoretical distribution through the transfer function to produce the “theoretically observable” distribution
- c-Perform Montecarlo realizations of the MBH population
- d-Create catalogs of observed binaries including FIM errors

We compare:

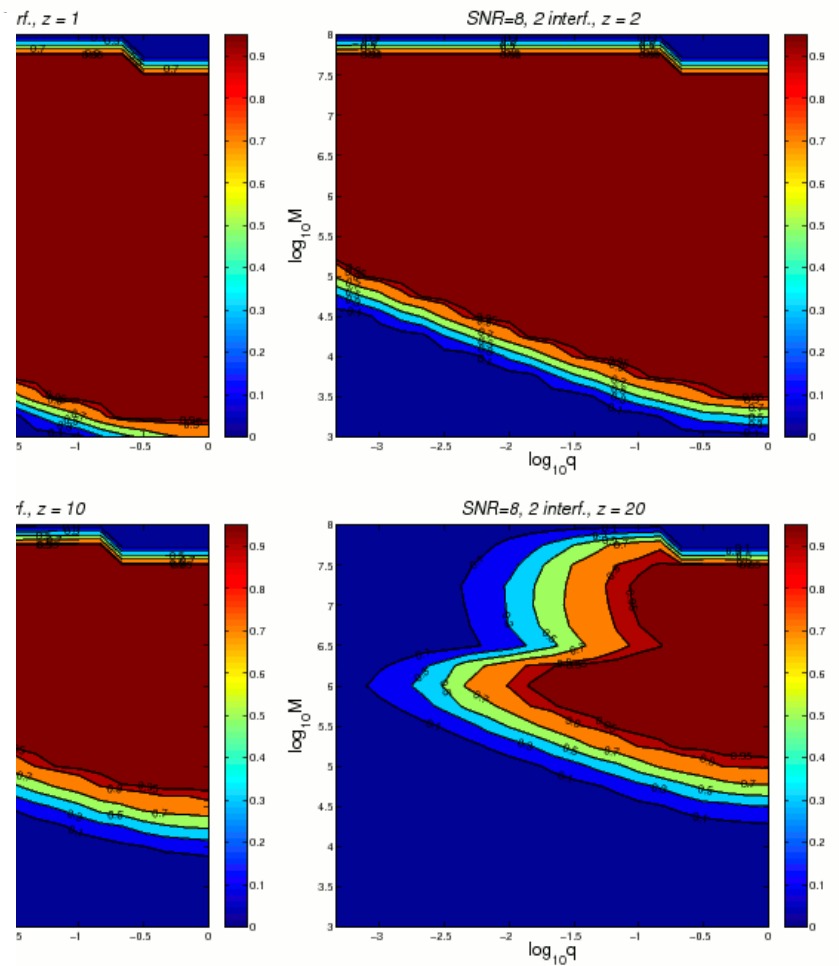
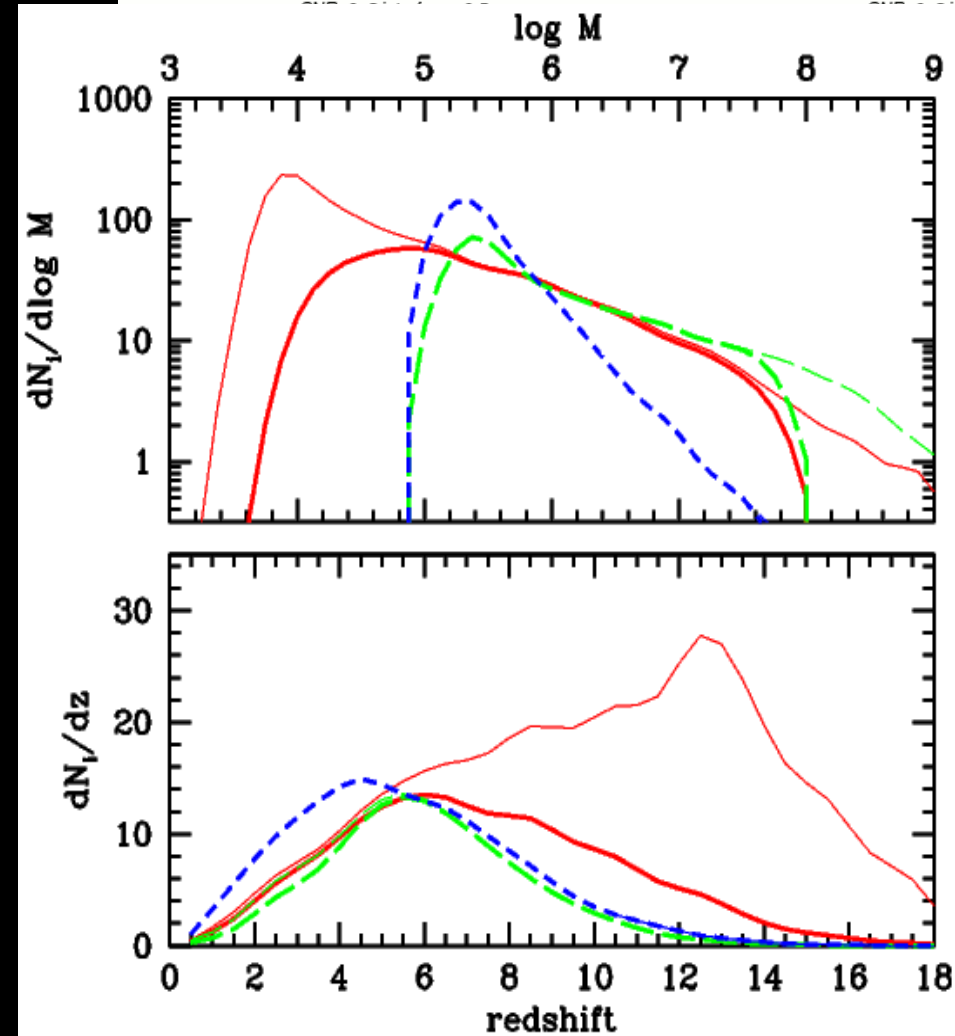
- the 10 pure models described before through the odds ratio
- artificially mixed models of the form $N_{mix} = f_1 N_1 + \dots + f_n N_n$
we find the maximum of the posterior distribution in the mixing parameter space
- consistently mixed models, again finding the maximum of the Posterior distribution

We consider the distribution $d^3N/dMdqdz$, we ignore spins.

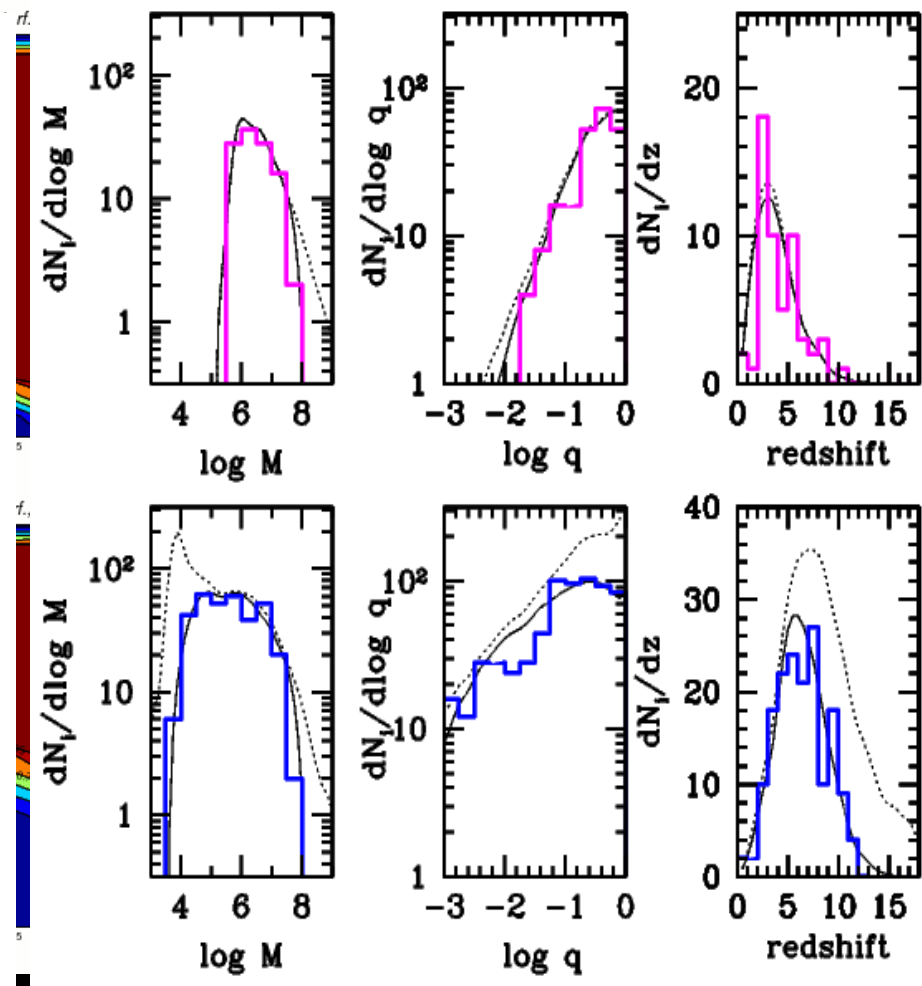
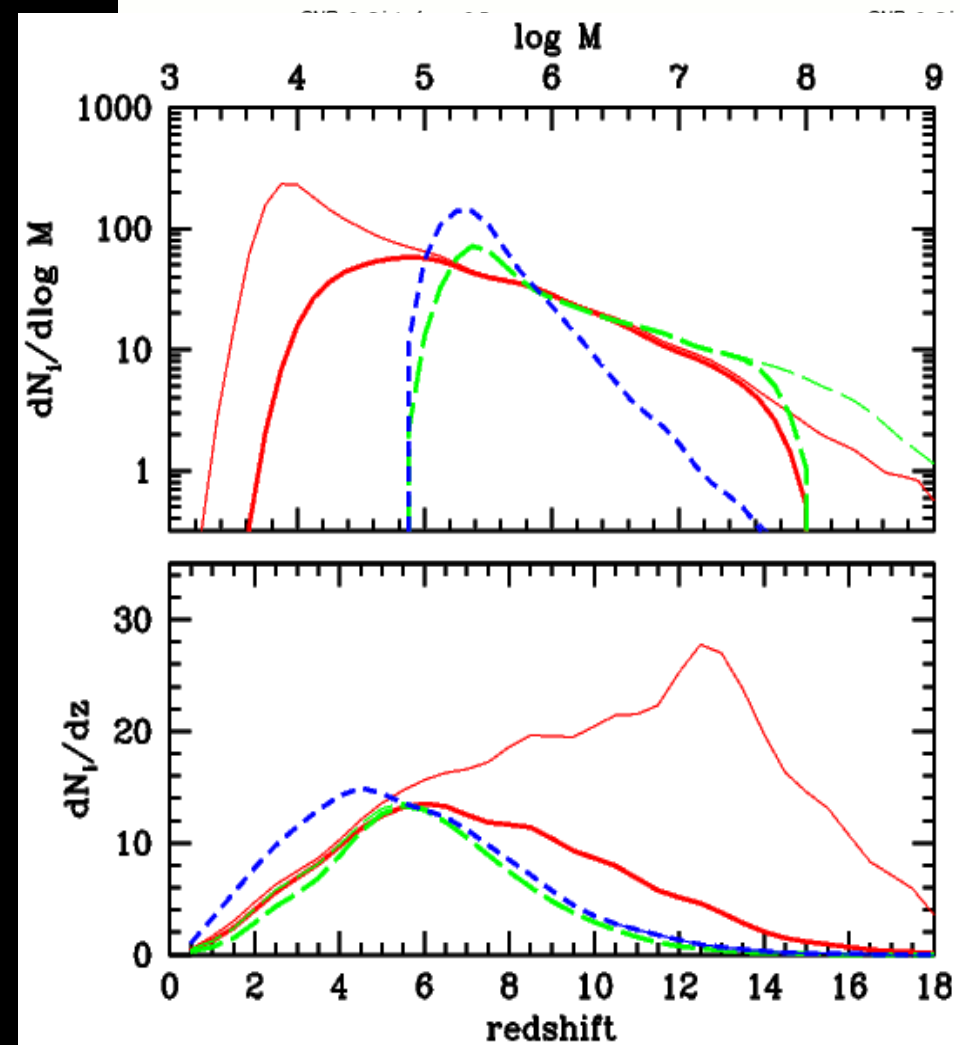
We use 2PN circular binary waveforms, no merger, no ringdown



$$N_{T_{i,j}}(z, M, q) = \frac{d^3 N_i}{dz dM dq} \times T_j(z, M, q)$$



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Likelihood of the dataset for a given choice of the parameters

$$p(D|\vec{\lambda}, M) = \prod_{i=1}^K \frac{(r_i(\vec{\lambda}))^{n_i} e^{-r_i(\vec{\lambda})}}{n_i!}$$

PURE MODEL COMPARISON: compute the odds ratio

$$O_{AB} = \frac{p(D|A) P(A)}{p(D|B) P(B)}$$

We assign confidence

$p_A = p(D|A)/(p(D|A)+p(D|B))$ to model A

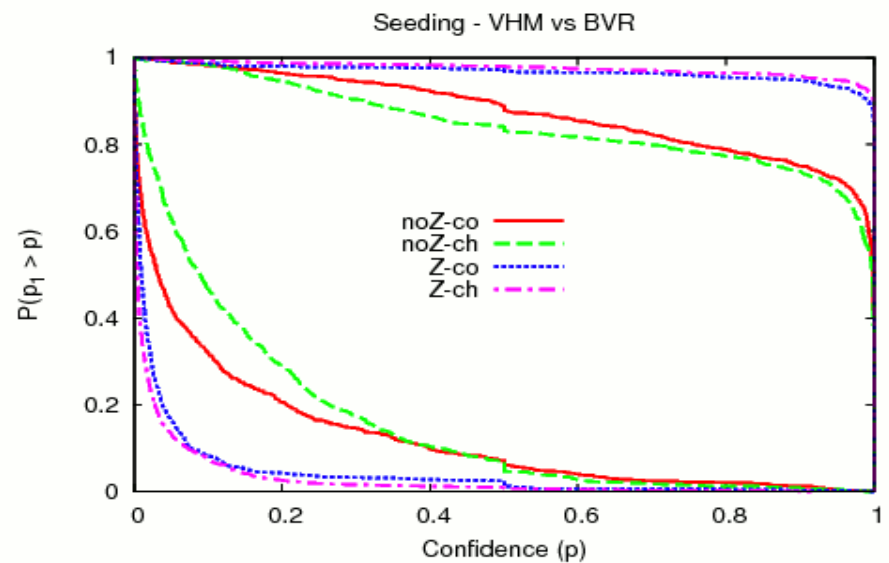
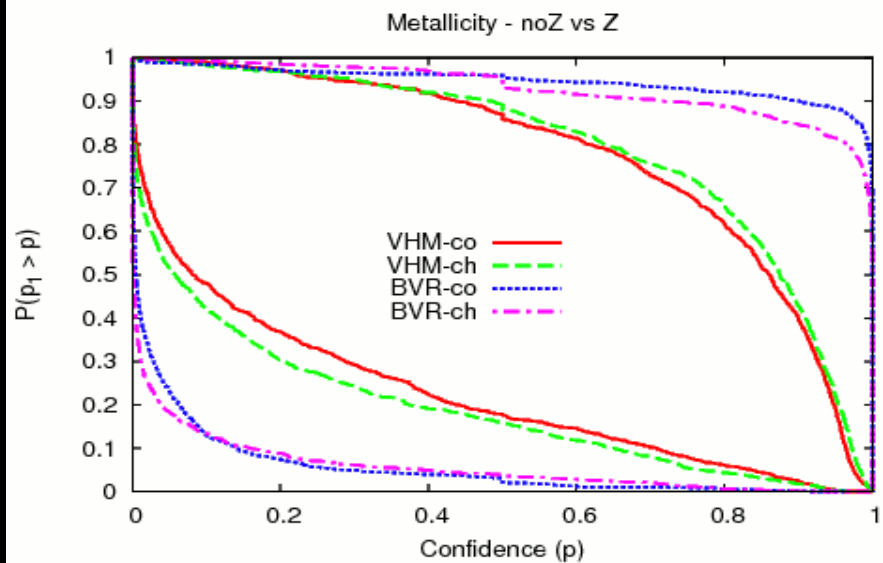
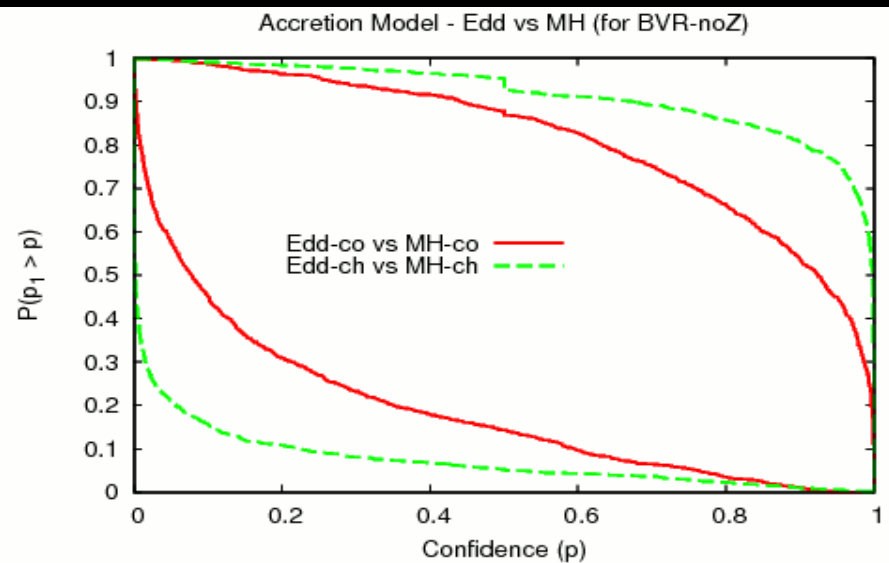
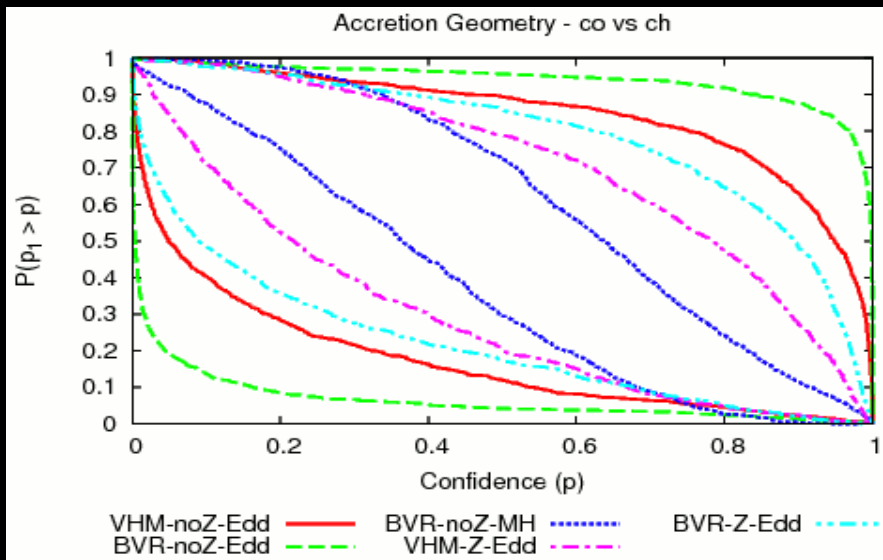
$p_B = 1 - p_A$ to model B

We build probability CDF over 1000 MC realization of the observed population

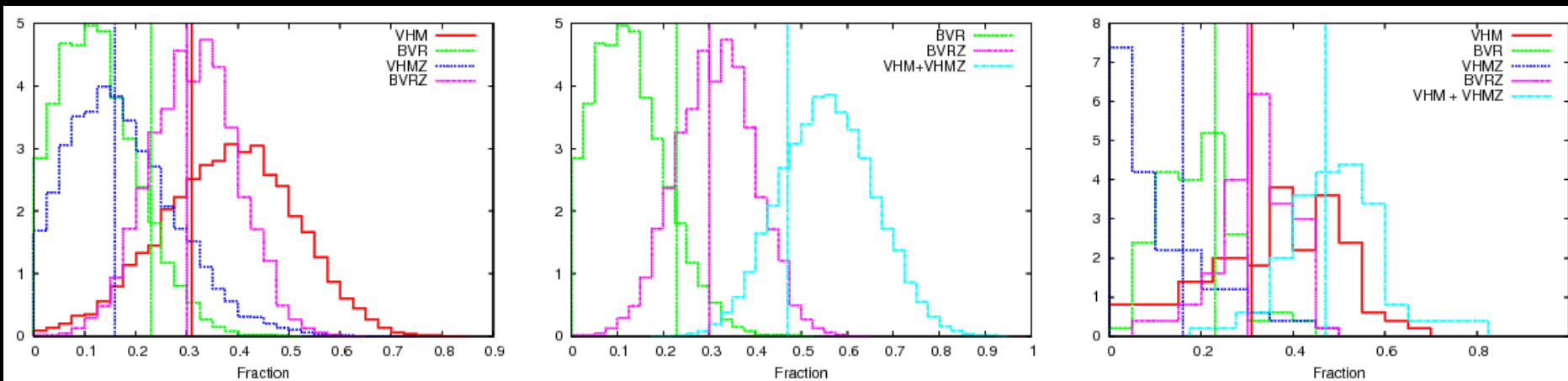
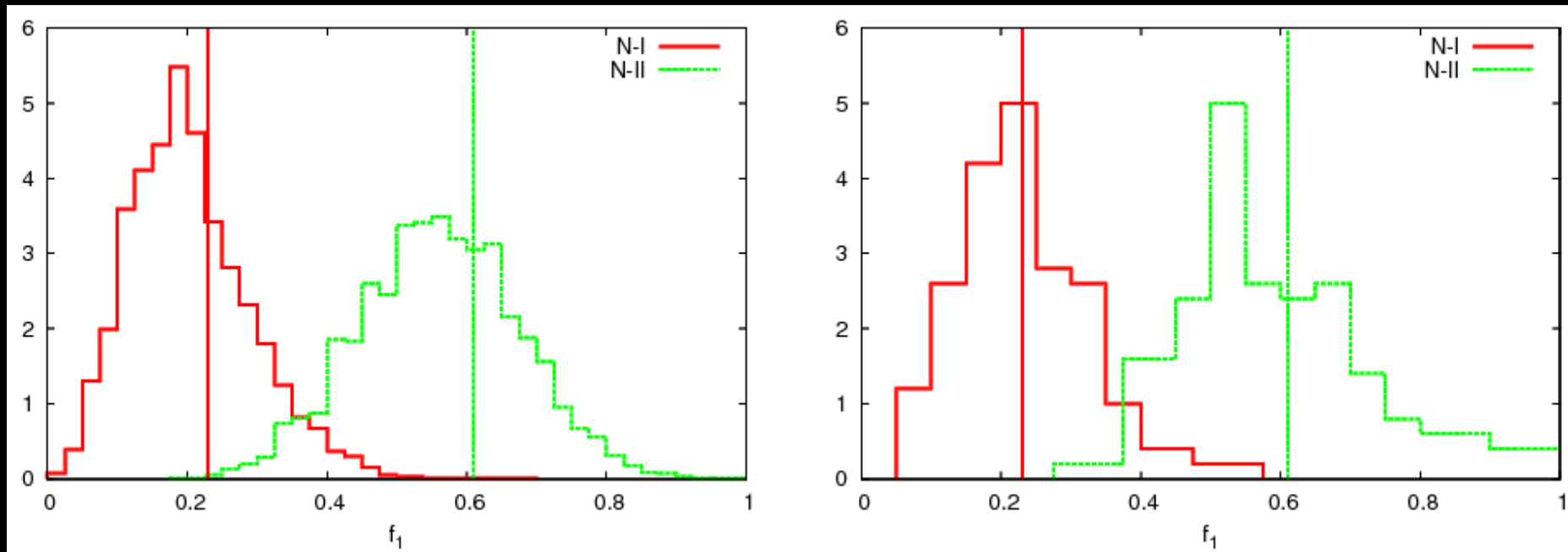
MIXED MODELS: find the maximum of the posterior

$$p(\vec{\lambda}|D, M) = \frac{p(D|\vec{\lambda}, M) \pi(\vec{\lambda})}{\mathcal{Z}},$$
$$\mathcal{Z} = \int p(D|\vec{\lambda}, M) \pi(\vec{\lambda}) d^N \lambda$$

Pure models: Cumulative distribution function for the confidence



Mixed models: evaluation of the posterior probability distribution function



CONCLUSIONS:

- LISA is a powerful instrument to assess MBH cosmic history**
- assuming 3yr observation we can virtually distinguish any pair of pure models at a 95% confidence level or more**
- when we mix different pure models, the mixing parameters can be recovered with 0.05-0.1 accuracy**
- we can also 'express' consistently mixed models as a combination of pure models.**

